



EQUILIBRIUM APPROACH IN THE DERIVATION OF DIFFERENTIAL EQUATIONS FOR HOMOGENEOUS ISOTROPIC MINDLIN PLATES

C. C. Ike*

DEPT OF CIVIL ENGINEERING, ENUGU STATE UNIVERSITY OF SCIENCE & TECHNOLOGY ENUGU, ENUGU STATE, NIGERIA

E-mail address: ikecc2007@yahoo.com

ABSTRACT

In this paper, the differential equations of Mindlin plates are derived from basic principles by simultaneous satisfaction of the differential equations of equilibrium, the stress-strain laws and the strain-displacement relations for isotropic, homogenous linear elastic materials. Equilibrium method was adopted in the derivation. The Mindlin plate equation was obtained as a system of simultaneous partial differential equations in terms of three displacement variables (parameters) namely $w(x, y, z = 0)$, $\theta_x(x, y)$ and $\theta_y(x, y)$ where $w(x, y, z)$ is the transverse displacement θ_x and θ_y are rotations of the middle surface. It was shown that when $\theta_x = -\frac{\partial w}{\partial x}$, $\theta_y = -\frac{\partial w}{\partial y}$ and $k \rightarrow \infty$, where k is the shear correction factor, the Mindlin plate equations reduce to the classical Kirchhoff plate equation which is a biharmonic equation in terms of $w(x, y, z = 0)$.

Keywords: *Mindlin plate, Kirchhoff plate, transverse displacement, rotations, shear correction factor, biharmonic equation.*

1. INTRODUCTION/LITERATURE REVIEW

Plates are three dimensional structural members frequently used as fluid containers (circular, elliptical and rectangular plates), building and bridge decks (rectangular and skewed slabs), aircraft wing panels (skewed plates), retaining walls, aerospace panels, and machine components [1 – 4]. Plates can be subjected to in plane loads and transverse loads, and can be simply supported, clamped or free at the edges.

The plate problem belongs to elasticity theory, and is usually to find the distribution of stress fields, strain fields and displacement fields in a given plate under known loading and support conditions. The exact solution is governed by a system of fifteen partial differential equations of equilibrium, material constitutive laws and kinematic equations that are solved subject to the loading conditions and the boundary conditions [5-8].

The basic idea of most plate theories is to reduce the three dimensional problem of plates to two dimensional problems. This reduction of dimension is accomplished by integrating out one of the dimensions, usually the plate's thickness, z , and thus expressing internal stresses using internal stress resultants, M_{xx} , M_{xy} , M_{yy} , Q_x and Q_y . All plate theories assume a kinematic assumption that the strains can be expanded in the smallest dimension [9, 8, 10]. Accordingly several theories had been formulated to describe the plate problems. Some of these plate

theories are Kirchhoff plate theory [3], Mindlin plate theory [4], Reissner plate theory [11, 12], Reddy plate theory [8], Levinson plate theory [13].

In this paper, the Mindlin plate theory for isotropic plates is formulated and derived using the equilibrium method, and presented as a problem of the theory of elasticity.

2. RESEARCH AIMS AND OBJECTIVES

The general aim and objective of this study is to derive the Mindlin plate theory from fundamental principles using the equilibrium approach. The specific aims and objectives are

- (i) to derive the governing equations of the Mindlin plate theory from fundamental principles of elasticity theory, using the equilibrium approach
- (ii) to show the relationship between the Mindlin plate theory and the Kirchhoff-Love plate theory.

3. DERIVATION OF MINDLIN PLATE EQUATIONS

The fundamental assumptions are as follows: [3, 4, 8, 14]

- (i) the state of deformation is described by the transverse displacement in the z -direction of the middle surface $w(x, y, z = 0)$ and the rotations θ_x and θ_y of the middle surface, where θ_x and θ_y are rotations about the x and y axes of lines normal to the middle surface before deformation [4].

- (ii) plane cross-sections originally perpendicular to the middle plane of the plate remain plane, but not necessarily orthogonal to the middle surface. Hence according to [14], $\varepsilon_r \neq 0$, $\varepsilon_{xz} \neq 0$.
- (iii) the middle surface remains neutral during bending, and is the neutral surface
- (iv) the displacement components vary linearly across the thickness
- (v) the plate material is isotropic, homogeneous and linear elastic

Let the displacement field be given by

$$u(x, y, z) = z\theta_x(x, y) \quad (1)$$

$$v(x, y, z) = z\theta_y(x, y) \quad (2)$$

$$w(x, y, z) = w(x, y) \quad (3)$$

In (1) to (3), u , v , w are displacement components in the x , y , z coordinate directions; and θ_x and θ_y are rotations about the x and y axis of lines normal to the middle surface before deformation. The kinematic (strain-displacement) equations for finite strain (small-displacement) elasticity problems are thus Equations (4) - (9):

$$\varepsilon_{xx} = \frac{\partial u}{\partial x} \quad (4)$$

$$\varepsilon_{yy} = \frac{\partial v}{\partial y} \quad (5)$$

$$\varepsilon_{zz} = \frac{\partial w}{\partial z} \quad (6)$$

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \quad (7)$$

$$\gamma_{xz} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \quad (8)$$

$$\gamma_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} \quad (9)$$

Where ε_{xx} , ε_{yy} , ε_{zz} are normal strains and γ_{xy} , γ_{yz} , γ_{xz} are shear strains. Substitution of Equations (1) - (3) into (4) - (9) gives:

$$\varepsilon_{xx} = z \frac{\partial \theta_x}{\partial x} \quad (10)$$

$$\varepsilon_{yy} = z \frac{\partial \theta_y}{\partial y} \quad (11)$$

$$\varepsilon_{zz} = 0 \quad (12)$$

$$\gamma_{xy} = z \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \quad (13)$$

$$\gamma_{xz} = \theta_x + \frac{\partial w}{\partial x} \quad (14)$$

$$\gamma_{yz} = \theta_y + \frac{\partial w}{\partial y} \quad (15)$$

Plate problems are required to satisfy the stress stain laws, and for two dimensional plane stress problems, we have

$$\sigma_{xx} = \frac{E}{1 - \mu^2} (\varepsilon_{xx} + \mu \varepsilon_{yy}) \quad (16)$$

$$\sigma_{yy} = \frac{E}{1 - \mu^2} (\varepsilon_{yy} + \mu \varepsilon_{xx}) \quad (17)$$

$$\sigma_{zz} = \mu (\sigma_{xx} + \sigma_{yy}) \quad \left(\text{NB: } \varepsilon_{yy} = 0 \right) \\ = \frac{1}{E} (\sigma_{zz} - \mu \sigma_{xx} - \mu \sigma_{yy}) \quad (18)$$

$$\tau_{xz} = \frac{E}{2(1 + \mu)} \gamma_{xz} \quad (19)$$

$$\tau_{xy} = \frac{E}{2(1 + \mu)} \gamma_{xy} \quad (20)$$

$$\tau_{yz} = \frac{E}{2(1 + \mu)} \gamma_{yz} \quad (21)$$

Here σ_{xx} , σ_{yy} are normal stresses, τ_{xy} is the shear stress. These stress-stain laws are expressed in terms of the displacement components by substituting Equations (10) - (15) into Equations (16) - (21) to obtain Equations (22) - (26).

$$\sigma_{xx} = \frac{Ez}{1 - \mu^2} \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) \quad (22)$$

$$\sigma_{yy} = \frac{Ez}{1 - \mu^2} \left(\frac{\partial \theta_y}{\partial y} + \mu \frac{\partial \theta_x}{\partial x} \right) \quad (23)$$

$$\tau_{xy} = \frac{Ez}{2(1 + \mu)} \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \quad (24)$$

$$\tau_{xz} = \frac{E}{2(1 + \mu)} \left(\theta_x + \frac{\partial w}{\partial x} \right) \quad (25)$$

$$\tau_{yz} = \frac{E}{2(1 + \mu)} \left(\theta_y + \frac{\partial w}{\partial y} \right) \quad (26)$$

We wish to eliminate the thickness variable z from the equations, by expressing the internal stresses σ_{xx} , σ_{yy} , σ_{zz} , τ_{xz} , τ_{yz} in terms of stress resultants or forces defined as M_{xx} , M_{yy} , M_{xy} , Q_x and Q_y where M_{xx} , M_{yy} are bending moments, M_{xy} is the twisting moment and Q_x and Q_y are shear forces.

From statics, the internal stress resultants are obtained as:

$$M_{xx} = \int_{-h/2}^{h/2} \sigma_{xx} z dz \quad (27)$$

$$M_{yy} = \int_{-h/2}^{h/2} \sigma_{yy} z dz \quad (28)$$

$$M_{xy} = \int_{-h/2}^{h/2} \tau_{xy} z dz \quad (29)$$

$$Q_x = \int_{-h/2}^{h/2} \tau_{xz} dz \quad (30)$$

$$Q_y = \int_{-h/2}^{h/2} \tau_{yz} dz \quad (31)$$

where $\frac{h}{2} \leq z \leq \frac{h}{2}$ and h is the plate thickness.

Using Equations (22) - (26), the internal stress resultants become

$$M_{xx} = \int_{-h/2}^{h/2} \frac{Ez}{1-\mu^2} \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) z dz \quad (32)$$

$$M_{xx} = \frac{Eh^3}{12(1-\mu^2)} \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) = D \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) \quad (33)$$

$$M_{yy} = D \left(\frac{\partial \theta_y}{\partial y} + \mu \frac{\partial \theta_x}{\partial x} \right) \quad (34)$$

$$M_{xy} = \frac{D(1-\mu)}{2} \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \quad (35)$$

$$Q_x = khG \left(\theta_x + \frac{\partial w}{\partial x} \right) \quad (36)$$

$$Q_y = khG \left(\theta_y + \frac{\partial w}{\partial y} \right) \quad (37)$$

where k is the shear correction factor

From the differential equations of equilibrium of an infinitesimal plate element,

$$\frac{\partial M_{xx}}{\partial x} + \frac{\partial M_{xy}}{\partial y} = Q_x \quad (38)$$

$$\frac{\partial M_{xy}}{\partial x} + \frac{\partial M_{yy}}{\partial y} = Q_y \quad (39)$$

$$\frac{\partial Q_x}{\partial x} + \frac{\partial Q_y}{\partial y} = -q \quad (40)$$

Substitution of Equations (32) - (37) into Equations (38) - (40), we obtain:

$$\begin{aligned} & \frac{\partial}{\partial x} \left(D \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) \right) \\ & + \frac{\partial}{\partial y} \left(\frac{D(1-\mu)}{2} \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \right) \\ & = khG \left(\theta_x + \frac{\partial w}{\partial x} \right) \end{aligned} \quad (41)$$

$$\begin{aligned} & \frac{\partial}{\partial x} \left(\frac{(1-\mu)}{2} D \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \right) \\ & + \frac{\partial}{\partial x} \left(D \left(\frac{\partial \theta_y}{\partial y} + \mu \frac{\partial \theta_x}{\partial x} \right) \right) \\ & = khG \left(\theta_y + \frac{\partial w}{\partial y} \right) \end{aligned} \quad (42)$$

$$\begin{aligned} & \frac{\partial}{\partial x} \left(khG \left(\theta_x + \frac{\partial w}{\partial x} \right) \right) + \frac{\partial}{\partial y} \left(khG \left(\theta_y + \frac{\partial w}{\partial y} \right) \right) \\ & = -q \end{aligned} \quad (43)$$

Simplification yields:

$$\nabla^2 w + \frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} = \frac{q}{kGh} \quad (44)$$

$$\begin{aligned} & D \left(\frac{\partial^2 \theta_x}{\partial x^2} + \mu \frac{\partial^2 \theta_y}{\partial x \partial y} \right) + \frac{(1-\mu)D}{2} \left(\frac{\partial^2 \theta_x}{\partial y^2} + \mu \frac{\partial^2 \theta_y}{\partial y \partial x} \right) \\ & = khG \left(\theta_x + \frac{\partial w}{\partial x} \right) \end{aligned} \quad (45)$$

$$\begin{aligned} & D \left(\frac{\partial^2 \theta_y}{\partial y^2} + \mu \frac{\partial^2 \theta_x}{\partial x \partial y} \right) + \frac{(1-\mu)D}{2} \left(\frac{\partial^2 \theta_x}{\partial y \partial x} + \mu \frac{\partial^2 \theta_y}{\partial x^2} \right) \\ & = khG \left(\theta_y + \frac{\partial w}{\partial y} \right) \end{aligned} \quad (46)$$

Equations (39) - (40) could be combined to obtain one partial differential equation of equilibrium as follows:

$$\frac{\partial^2 M_{xx}}{\partial x^2} + 2 \frac{\partial^2 M_{xy}}{\partial x \partial y} + \frac{\partial^2 M_{yy}}{\partial y^2} = -q \quad (47)$$

The substitution of Equations (33) - (35) in Equation (47) would yield

$$\begin{aligned} & \frac{D \partial^2}{\partial x^2} \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) + \mu \frac{D(1-\mu)}{2} \frac{\partial^2}{\partial x \partial y} \left(\frac{\partial \theta_x}{\partial y} + \frac{\partial \theta_y}{\partial x} \right) \\ & + \frac{D \partial^2}{\partial y^2} \left(\frac{\partial \theta_y}{\partial y} + \mu \frac{\partial \theta_x}{\partial x} \right) = -q \end{aligned} \quad (48)$$

Simplification yields:

$$\nabla^2 \left(\frac{\partial \theta_x}{\partial x} + \mu \frac{\partial \theta_y}{\partial y} \right) = -\frac{q}{D} \quad (49)$$

where ∇^2 is the Laplacian operator.

The governing equation of isotropic homogeneous Mindlin plates are thus obtained as a system of three unknown displacement parameters w , θ_x and θ_y . They are Equations (44), (45), (46) and (49).

4. RELATIONSHIP BETWEEN THE MINDLIN PLATE THEORY AND THE CLASSICAL KIRCHHOFF PLATE THEORY

In the classical Kirchhoff plate theory, the shear strains γ_{xz} and γ_{yz} are disregarded; and assumed to be equal to zero, respectively. Then Equations (14) and (15) would yield

$$\gamma_{xz} = \theta_x + \frac{\partial w}{\partial x} = 0 \quad (50)$$

$$\gamma_{yz} = \theta_y + \frac{\partial w}{\partial y} = 0 \quad (51)$$

Hence,

$$\theta_x = -\frac{\partial w}{\partial x} \quad (52)$$

$$\theta_y = -\frac{\partial w}{\partial y} \quad (53)$$

The internal bending moment resultants Equations (33) - (35) would then become for the Kirchhoff plate theory:

$$M_{xx} = D \left(\frac{\partial}{\partial x} \left(-\frac{\partial w}{\partial x} \right) + \mu \frac{\partial}{\partial y} \left(-\frac{\partial w}{\partial y} \right) \right) \quad (54)$$

$$M_{xx} = -D \left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} \right) \quad (55)$$

$$M_{yy} = -D \left(\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \right) \quad (56)$$

$$M_{xy} = -D(1-\mu) \frac{\partial^2 w}{\partial x \partial y} \quad (57)$$

The shear force resultants are

$$Q_x = 0 \quad (58)$$

$$Q_y = 0 \quad (59)$$

The equation of equilibrium of Mindlin plate when shear strains are disregarded becomes:

$$\begin{aligned} \nabla^2 \left(\frac{\partial \theta_x}{\partial x} + \frac{\partial \theta_y}{\partial y} \right) &= \nabla^2 \left(\frac{\partial}{\partial x} \left(-\frac{\partial w}{\partial x} \right) + \frac{\partial}{\partial y} \left(-\frac{\partial w}{\partial x} \right) \right) \\ &= -\frac{q}{D} \end{aligned} \quad (60)$$

$$-\nabla^2 \nabla^2 w = -\frac{q}{D} \quad \text{or} \quad \nabla^2 \nabla^2 w = \frac{q}{D} = \nabla^4 w \quad (61)$$

where ∇^2 is the Laplacian operator while $\nabla^4 = \nabla^2 \nabla^2$ is the biharmonic operator.

Equation (61) is the governing partial differential equation of equilibrium for Kirchhoff plates. The Mindlin plate governing equation thus reduces to the classical Kirchhoff plate equation when $\varphi_x = \frac{\partial w}{\partial x}$ and $\varphi_y = -\frac{\partial w}{\partial y}$

5. DISCUSSIONS AND CONCLUSIONS

The Mindlin plate theory has been derived for isotropic, homogeneous linear elastic plates using the equilibrium method, taking into consideration the shear strain, and hence the shear stress across the plate thickness. However, the shear strain is assumed to be constant across the plate thickness, violating the predictions of the theory of elasticity since the shear stress is known to vary parabolically over the plate thickness. A shear correction factor k , is introduced to ensure that the correct amount of internal energy is predicted by the Mindlin plate theory. The shear correction factor merely yields a resultant shear stress that agrees with the predictions of elasticity theory, but the distribution of shear stress and shear strain violate the theory of elasticity solutions. The transverse shear stresses are found to be constant over the plate thickness, contradicting the shear free boundary condition on the plate surfaces.

The differential equations for Mindlin plates have been derived in this paper from fundamental principles using the equilibrium method. The equations were derived using the requirements of the differential equations of equilibrium for an infinitesimal element of the plate, the material constitutive laws, and the strain-displacement relations for small-displacement elasticity problems. This three sets of relations were assumed to be satisfied simultaneously for any differential element of the plate. The deformation field was assumed, in line with Mindlin's plate to be defined by three displacement components, namely the transverse displacement w , and the rotations θ_x and θ_y relaxing the orthogonality requirements of plane cross sections to the middle (neutral) surface, by enforcing $\varepsilon_{xy} \neq 0$ and $\varepsilon_{xz} \neq 0$ shear strains were accounted for in the derivation, yielding stress-strain laws given by Equations (16) – (21). The material constitutive laws expressed in terms of the

three displacement fields were obtained as Equations (22) – (26). The internal stress resultants were obtained as Equations (33) – (37). Differential equations of equilibrium expressed in terms of the internal stress resultants were then used to obtain the governing partial differential equations of equilibrium of Mindlin plates as Equations (44) – (46) and (49). The governing equations of Mindlin plates are a system of partial differential equations in terms of three unknown displacement parameters, namely w , θ_x and θ_y .

Mindlin theory is a two dimensional plate theory and cannot yield an exact solution to plate problems since plate problems are three dimensional. The major merits of the Mindlin plate theory are that the theory takes consideration of shear effects, and the theory simplifies to the classical Kirchhoff plate theory by setting $\theta_x = -\frac{\partial w}{\partial x}$, $\theta_y = -\frac{\partial w}{\partial y}$, and $k \rightarrow \infty$. The shear correction factor introduced to account for the non-uniformity of the shear strain on the cross-section is dependent on the shape of the cross-section, and the shear stress distribution. Reissner suggests that $k=5/6$.

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